

The Hermite polynomials build an orthogonal system. This paper checks the orthogonality. It is based on:

“Hort-Thoma, Die Differentialgleichungen der Technik und Physik”, Leipzig 1950,

Griffiths, Introduction to quantum mechanics, third edition, page 54.

https://www.ruf.rice.edu/~baring/phys516/phys516_2021_lec_041321.pdf

For the proof of orthogonality:

<https://math.stackexchange.com/questions/2701650/orthogonality-and-norm-of-hermite-polynomials>

Hope I can help you with learning quantum mechanics.

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Hermite polynomials

We begin with Griffiths Problem 2.16 to become acquainted with Hermite polynomials.

We use Rodrigues formula:

$$H_n(\xi) = (-1)^n \cdot e^{\xi^2} \cdot \left(\left(\frac{d}{d\xi} \right)^n e^{-\xi^2} \right) \quad (1)$$

We can invert this:

$$\left(\frac{d}{d\xi} \right)^n e^{-\xi^2} = H_n(\xi) \cdot (-1)^n \cdot e^{-\xi^2} \quad (1a)$$

Note: Instead of x we use $\xi = \sqrt{\frac{m\omega}{\hbar}} \cdot x$, a dimensionless variable.

We derive $H_0(\xi), H_1(\xi), H_2(\xi), H_3(\xi), H_4(\xi)$

$$H_0(\xi) = 1$$

$$H_1(\xi) = 2 \cdot \xi$$

$$H_2(\xi) = (-1)^2 e^{\xi^2} \left(\frac{d}{d\xi} \right)^2 e^{-\xi^2} = e^{\xi^2} \frac{d}{d\xi} \left(\frac{d}{d\xi} e^{-\xi^2} \right) =$$

$$e^{\xi^2} \frac{d}{d\xi} (-2 \cdot \xi \cdot e^{-\xi^2}) = e^{\xi^2} (-2 \cdot e^{-\xi^2} + 4 \cdot \xi^2 \cdot e^{-\xi^2}) =$$

$$e^{\xi^2} (-2 \cdot e^{-\xi^2} + 4 \cdot \xi^2 \cdot e^{-\xi^2}) = -2 + 4 \cdot \xi^2$$

$$H_3(\xi) = (-1)^3 e^{\xi^2} \left(\frac{d}{d\xi} \right)^3 e^{-\xi^2} = -e^{\xi^2} \frac{d}{d\xi} \left(\frac{d}{d\xi} \left(\frac{d}{d\xi} e^{-\xi^2} \right) \right) =$$

$$-e^{\xi^2} \frac{d}{d\xi} (-2 \cdot e^{-\xi^2} + 4 \cdot \xi^2 \cdot e^{-\xi^2}) = -e^{\xi^2} (4 \cdot \xi \cdot e^{-\xi^2} + 8 \cdot \xi \cdot e^{-\xi^2} - 8 \cdot \xi^3 \cdot e^{-\xi^2}) =$$

$$8 \cdot \xi^3 - 12 \cdot \xi$$

$$H_4(\xi) = (-1)^4 e^{\xi^2} \left(\frac{d}{d\xi} \right)^4 e^{-\xi^2} = e^{\xi^2} \frac{d}{d\xi} \left(\frac{d}{d\xi} \left(\frac{d}{d\xi} \left(\frac{d}{d\xi} e^{-\xi^2} \right) \right) \right) =$$

$$e^{\xi^2} \frac{d}{d\xi} (12 \cdot \xi \cdot e^{-\xi^2} - 8 \cdot \xi^3 \cdot e^{-\xi^2}) =$$

$$e^{\xi^2} (12 \cdot e^{-\xi^2} - 24 \cdot \xi^2 \cdot e^{-\xi^2} - 24 \cdot \xi^2 \cdot e^{-\xi^2} + 16 \cdot \xi^4 \cdot e^{-\xi^2}) =$$

$$16 \cdot \xi^4 - 48 \cdot \xi^2 + 12$$

This should be enough to trust in the Rodrigues formula.

We have the recursion relation:

$$H_{n+2}(\xi) = 2 \cdot \xi \cdot H_{n+1}(\xi) - 2 \cdot (n+1) \cdot H_n(\xi)$$

With the recursion relation we build $H_5(\xi)$ and $H_6(\xi)$:

$$\begin{aligned}
 H_5(\xi) &= 2 \cdot \xi \cdot H_4(\xi) - 2 \cdot 4 \cdot H_3(\xi) = \\
 &2 \cdot \xi \cdot (16 \cdot \xi^4 - 48 \cdot \xi^2 + 12) - 8 \cdot (8 \cdot \xi^3 - 12 \cdot \xi) = \\
 &32 \cdot \xi^5 - 96 \cdot \xi^3 + 24 \cdot \xi - 64 \cdot \xi^3 + 96 \cdot \xi = \\
 &32 \cdot \xi^5 - 160 \cdot \xi^3 + 120 \cdot \xi \\
 H_6(\xi) &= 2 \cdot \xi \cdot H_5(\xi) - 2 \cdot 5 \cdot H_4(\xi) = \\
 &2 \cdot \xi \cdot (32 \cdot \xi^5 - 160 \cdot \xi^3 + 120 \cdot \xi) - 10 \cdot (16 \cdot \xi^4 - 48 \cdot \xi^2 + 12) = \\
 &64 \cdot \xi^6 - 320 \cdot \xi^4 + 240 \cdot \xi^2 - 160 \cdot \xi^4 + 480 \cdot \xi^2 - 120 = \\
 &64 \cdot \xi^6 - 480 \cdot \xi^4 + 720 \cdot \xi^2 - 120
 \end{aligned}$$

For Hermite polynomials holds:

$$\frac{d}{d\xi} H_n(\xi) = 2 \cdot n \cdot H_{n-1}(\xi)$$

We check this by differentiating $H_6(\xi)$ and $H_5(\xi)$

$$\begin{aligned}
 \frac{d}{d\xi} H_6(\xi) &= \frac{d}{d\xi} (64 \cdot \xi^6 - 480 \cdot \xi^4 + 720 \cdot \xi^2 - 120) = \\
 &384 \cdot \xi^5 - 1920 \cdot \xi^3 + 1440 \cdot \xi = \\
 &2 \cdot 6 \cdot (32 \cdot \xi^5 - 160 \cdot \xi^3 + 120 \cdot \xi) = \\
 &12 \cdot H_5(\xi) \\
 \frac{d}{d\xi} H_5(\xi) &= \frac{d}{d\xi} (32 \cdot \xi^5 - 160 \cdot \xi^3 + 120 \cdot \xi) = \\
 &160 \cdot \xi^4 - 480 \cdot \xi^2 + 120 = \\
 &2 \cdot 5 \cdot (16 \cdot \xi^4 - 48 \cdot \xi^2 + 12) = \\
 &10 \cdot H_4(\xi)
 \end{aligned}$$

The generating function for the Hermite polynomials:

$$e^{(-z^2+2 \cdot z \cdot \xi)}$$

We expand into Taylor power series at $z = 0$:

$$\begin{aligned}
 e^{(-z^2+2 \cdot z \cdot \xi)} &= \frac{d^0}{dz^0} e^{(-z^2+2 \cdot z \cdot \xi)} \cdot z^0 + \frac{d^1}{dz^1} e^{(-z^2+2 \cdot z \cdot \xi)} \cdot z^1 + \frac{d^2}{dz^2} e^{(-z^2+2 \cdot z \cdot \xi)} \cdot z^2 + \frac{d^3}{dz^3} e^{(-z^2+2 \cdot z \cdot \xi)} \cdot z^3 + \dots = \\
 &\sum_{n=0}^{\infty} \frac{d^n}{dz^n} e^{(-z^2+2 \cdot z \cdot \xi)} \cdot z^n
 \end{aligned}$$

We build $H_1(\xi), H_2(\xi), H_3(\xi)$.

$$H_1(\xi) = \frac{d^1}{dz^1} e^{(-z^2+2 \cdot z \cdot \xi)} =$$

$$(-2 \cdot z + 2 \cdot \xi) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} = (\text{note: } z = 0) = 2 \cdot \xi$$

$$H_2(\xi) = \frac{\frac{d^2}{dz^2} e^{(-z^2+2 \cdot z \cdot \xi)}}{2!} =$$

$$\frac{d}{dz} \left((-2 \cdot z + 2 \cdot \xi) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} \right) =$$

$$(-2) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} + (-2 \cdot z + 2 \cdot \xi) \cdot (-2 \cdot z + 2 \cdot \xi) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} =$$

$$(-2 + 4 \cdot z^2 - 8 \cdot z \cdot \xi + 4 \cdot \xi^2) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} = (\text{note: } z = 0) = -2 + 4 \cdot \xi^2$$

$$H_3(\xi) = \frac{\frac{d^3}{dz^3} e^{(-z^2+2 \cdot z \cdot \xi)}}{3!} =$$

$$\frac{d}{dz} \left((-2 + 4 \cdot z^2 - 8 \cdot z \cdot \xi + 4 \cdot \xi^2) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} \right) =$$

$$\left((8 \cdot z - 8 \cdot \xi) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} + (-2 + 4 \cdot z^2 - 8 \cdot z \cdot \xi + 4 \cdot \xi^2) \cdot (-2 \cdot z + 2 \cdot \xi) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} \right) =$$

$$\left((8 \cdot z - 8 \cdot \xi) + (-2 + 4 \cdot z^2 - 8 \cdot z \cdot \xi + 4 \cdot \xi^2) \cdot (-2 \cdot z + 2 \cdot \xi) \right) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} =$$

$$(8 \cdot z - 8 \cdot \xi + 4 \cdot z - 4 \cdot \xi - 8 \cdot z^3 + 8 \cdot z^2 \cdot \xi + 16 \cdot z^2 \cdot \xi - 16 \cdot z \cdot \xi^2 - 8 \cdot z \cdot \xi^2 + 8 \cdot \xi^3) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} =$$

$$(12 \cdot z - 12 \cdot \xi - 8 \cdot z^3 + 24 \cdot z^2 \cdot \xi - 24 \cdot z \cdot \xi^2 + 8 \cdot \xi^3) \cdot e^{(-z^2+2 \cdot z \cdot \xi)} = (\text{note: } z = 0) = (-12 \cdot \xi + 8 \cdot \xi^3)$$

This might be enough to let us trust in the generating function.

We summarize:

$$H_0(\xi) = 1$$

$$H_1(\xi) = 2 \cdot \xi$$

$$H_2(\xi) = 4 \cdot \xi^2 - 2$$

$$H_3(\xi) = 8 \cdot \xi^3 - 12 \cdot \xi$$

$$H_4(\xi) = 16 \cdot \xi^4 - 48 \cdot \xi^2 + 12$$

$$H_5(\xi) = 32 \cdot \xi^5 - 160 \cdot \xi^3 + 120 \cdot \xi$$

$$H_6(\xi) = 64 \cdot \xi^6 - 480 \cdot \xi^4 + 720 \cdot \xi^2 - 120$$

We note that the leading coefficient of each polynomial is 2^n . (2)

Orthogonality

The wave function for the quantum harmonic oscillator:

$$\psi_n(\xi) = e^{-\frac{\xi^2}{2}} \cdot H_n(\xi)$$

We have to proof for $m \neq n, m < n$:

$$\int_{-\infty}^{\infty} \psi_n(\xi) \cdot \psi_m(\xi) d\xi = 0$$

We use Rodriguez formula (1):

$$H_m(\xi) = (-1)^m e^{\xi^2} \left(\frac{d}{d\xi} \right)^m e^{-\xi^2}$$

We resolve the wave functions:

$$\begin{aligned} \int_{-\infty}^{\infty} e^{-\frac{\xi^2}{2}} \cdot H_m(\xi) \cdot e^{-\frac{\xi^2}{2}} \cdot H_n(\xi) d\xi &= \int_{-\infty}^{\infty} e^{-\xi^2} \cdot H_m(\xi) \cdot H_n(\xi) d\xi = \\ \int_{-\infty}^{\infty} e^{-\xi^2} \cdot H_m(\xi) \cdot \left((-1)^n \cdot e^{\xi^2} \cdot \left(\left(\frac{d}{d\xi} \right)^n e^{-\xi^2} \right) \right) d\xi &= \\ (-1)^n \int_{-\infty}^{\infty} H_m(\xi) \cdot \left(\left(\frac{d}{d\xi} \right)^n e^{-\xi^2} \right) d\xi &=; \end{aligned}$$

We omit the factor $(-1)^n$:

$$\int_{-\infty}^{\infty} H_m(\xi) \cdot \left(\left(\frac{d}{d\xi} \right)^n e^{-\xi^2} \right) d\xi$$

We use partial integration:

$$\int f \cdot g' = [f \cdot g] - \int f' \cdot g$$

We assign:

$$\begin{aligned} f &:= H_m(\xi) \rightarrow f^{(1)} = \frac{d}{d\xi} (H_m(\xi)) \\ g^{(1)} &:= \left(\frac{d}{d\xi} \right)^n e^{-\xi^2} \rightarrow g = \left(\frac{d}{d\xi} \right)^{(n-1)} e^{-\xi^2} \end{aligned}$$

We get for one step:

$$\begin{aligned} \int_{-\infty}^{\infty} H_m(\xi) \cdot \left(\left(\frac{d}{d\xi} \right)^n e^{-\xi^2} \right) d\xi \rightarrow \\ \left[H_m(\xi) \cdot \left(\left(\frac{d}{d\xi} \right)^{(n-1)} e^{-\xi^2} \right) \right]_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \left(\frac{d}{d\xi} \right)^{(1)} (H_m(\xi)) \cdot \left(\left(\frac{d}{d\xi} \right)^{(n-1)} e^{-\xi^2} \right) d\xi \end{aligned}$$

We repeat this n -times and get:

$$\left[H_m(\xi) \cdot \left(\frac{d}{d\xi} \right)^{(n-n)} e^{-\xi^2} \right]_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \left(\frac{d}{d\xi} \right)^n (H_m(\xi)) \cdot \left(\frac{d}{d\xi} \right)^{(n-n)} e^{-\xi^2} d\xi$$

$$[H_m(\xi) \cdot (e^{-\xi^2})]_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \left(\frac{d}{d\xi} \right)^n (H_m(\xi)) \cdot (e^{-\xi^2}) d\xi$$

We examine: Due to $e^{-\xi^2}$ we have:

$$[H_m(\xi) \cdot (e^{-\xi^2})]_{-\infty}^{\infty} = 0$$

The integral:

$$\int_{-\infty}^{\infty} \left(\frac{d}{d\xi} \right)^n (H_m(\xi)) \cdot (e^{-\xi^2}) d\xi$$

This integral is zero because $(H_m(\xi))$ is a polynomial of grade $m < n$, so the derivatives becomes zero beginning with $n = m + 1$.

Case $m = n$. We have the same process as above but this time ending with:

$$\left[H_n(\xi) \cdot \left(\frac{d}{d\xi} \right)^{(n-n)} e^{-\xi^2} \right]_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \left(\frac{d}{d\xi} \right)^n (H_n(\xi)) \cdot \left(\frac{d}{d\xi} \right)^{(n-n)} e^{-\xi^2} d\xi =$$

$$[H_n(\xi) \cdot (e^{-\xi^2})]_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \left(\frac{d}{d\xi} \right)^n (H_n(\xi)) \cdot (e^{-\xi^2}) d\xi =;$$

As before we have:

$$[H_n(\xi) \cdot (e^{-\xi^2})]_{-\infty}^{\infty} = 0$$

We look at:

$$\int_{-\infty}^{\infty} \left(\frac{d}{d\xi} \right)^n (H_n(\xi)) \cdot e^{-\xi^2} d\xi$$

We take a closer look at:

$$\left(\frac{d}{d\xi} \right)^n (H_n(\xi)) \tag{3}$$

$H_n(\xi)$ is a polynomial in ξ of grade ξ^n . Deriving this polynomial n -times gives the factor $n!$

The leading coefficient of $H_n(\xi)$ is 2^n as we saw above (2).

With this we get the value of the integral:

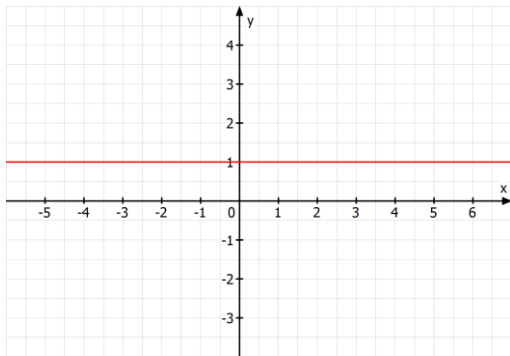
$$\int_{-\infty}^{\infty} \left(\frac{d}{d\xi} \right)^n (H_n(\xi)) \cdot e^{-\xi^2} d\xi = n! \cdot 2^n \cdot \int_{-\infty}^{\infty} e^{-\xi^2} d\xi =$$

$$n! \cdot 2^n \cdot \sqrt{\pi}$$

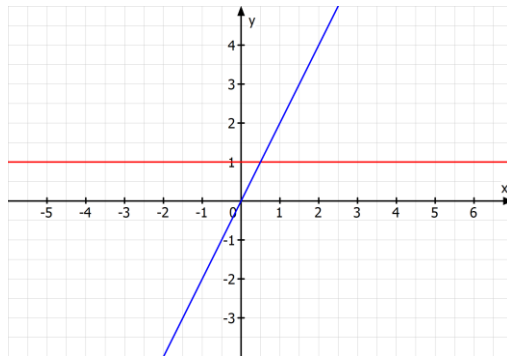
Note: If $m = n$ the Hermite polynomial disappears from the integral

Visualization

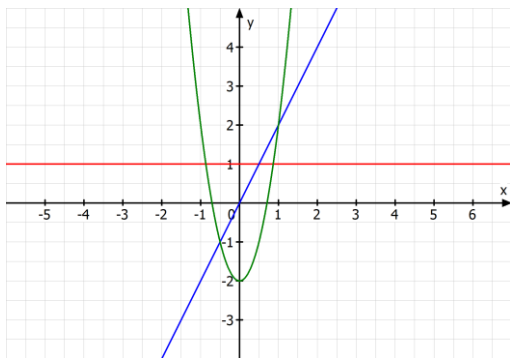
We plot the Hermite polynomials from $H_0(\xi)$ to $H_6(\xi)$.



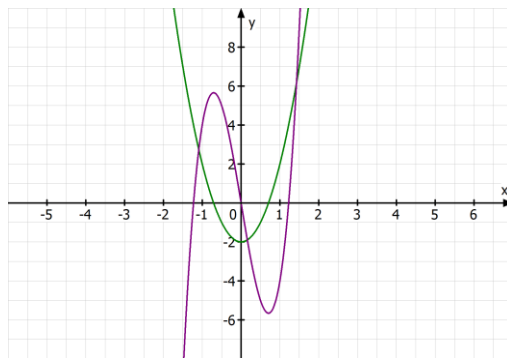
$H_0(\xi)$



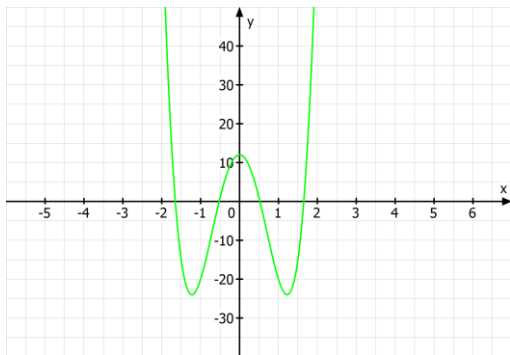
$H_1(\xi)$



$H_0(\xi) H_1(\xi) H_2(\xi)$



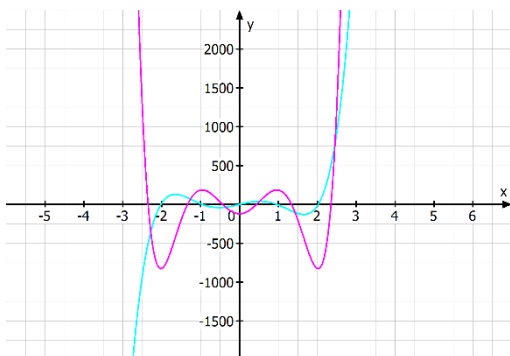
$H_2(\xi) H_3(\xi)$



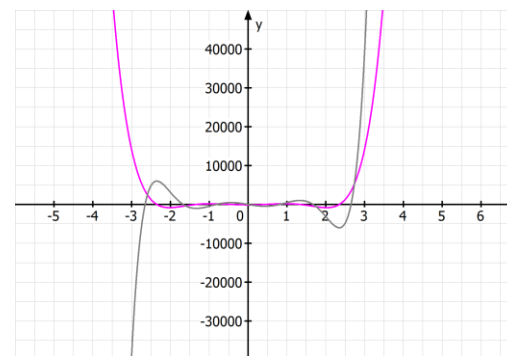
$H_4(\xi)$



$H_5(\xi)$



$H_5(\xi) H_6(\xi)$



$H_6(\xi) H_7(\xi)$

What we learn out of these plots is that the inner structure of the Hermite polynomials is interesting for small ξ only. If the value of ξ exceeds ± 3 , the Hermite polynomials reduce to the effect of the leading power of ξ .

We want to check what dimensions we are working with and use the definition of ξ :

$$\xi := \sqrt{\frac{m \cdot \omega}{\hbar}} \cdot x \rightarrow \sqrt{\frac{m \cdot 4 \cdot \pi^2 \cdot \nu}{h}} \cdot x$$

We need an example and choose a nonrelativistic electron.

Planck constant:	$h = 6.626 \cdot 10^{-34} \text{J} \cdot \text{s}$
Mass of an electron:	$m_0 = 9.109 \cdot 10^{-31} \text{kg}$
Speed of light:	$2.998 \cdot 10^8 \frac{\text{m}}{\text{s}}$
de Broglie relation for frequency:	$\nu = \frac{E}{h}$
From special relativity:	$E = \frac{m_0 \cdot c^2}{\sqrt{1 - \frac{v^2}{c^2}}}$

We assume $v \ll c$ and combine:

$$\begin{aligned} \xi &:= \sqrt{\frac{m_0 \cdot 4 \cdot \pi^2 \cdot \nu}{h}} \cdot x \rightarrow \sqrt{\frac{m_0 \cdot 4 \cdot \pi^2 \cdot E}{h^2}} \cdot x \rightarrow \sqrt{\frac{m_0 \cdot 4 \cdot \pi^2 \cdot m_0 \cdot c^2}{h^2}} \rightarrow \frac{m_0 \cdot 2 \cdot \pi \cdot c}{h} \\ \xi &= \frac{9.109 \cdot 10^{-31} \text{kg} \cdot 2 \cdot \pi \cdot 2.998 \cdot 10^8 \frac{\text{m}}{\text{s}}}{6.626 \cdot 10^{-34} \text{J} \cdot \text{s}} \cdot x \end{aligned}$$

We check the dimensions:

$$\left[\frac{\text{kg} \cdot \text{m}}{\text{J} \cdot \text{s}^2} \cdot \text{m} \right] = \left[\frac{\text{kg} \cdot \text{m} \cdot \text{s}^2}{\text{kg} \cdot \text{m}^2 \cdot \text{s}^2} \cdot \text{m} \right] = \left[\frac{1}{\text{m}} \cdot \text{m} \right] = 1$$

ξ is a dimensionless variable.

We get ξ for a slow-moving electron, $v \ll c$:

$$\xi \approx 2.590 \cdot 10^{12} \cdot x$$

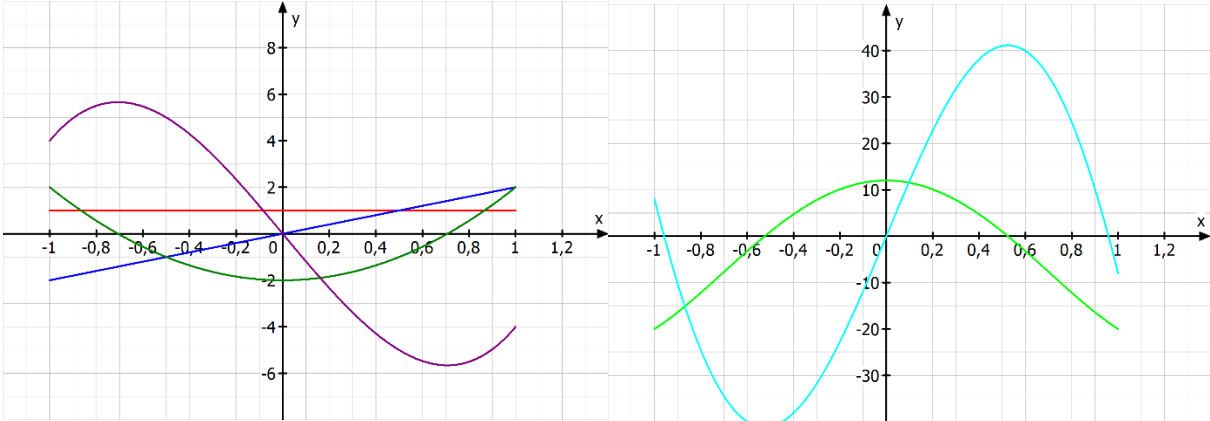
We multiply with the classical electron radius:

$$r_e \approx 2.818 \cdot 10^{-15} \text{m}$$

We get:

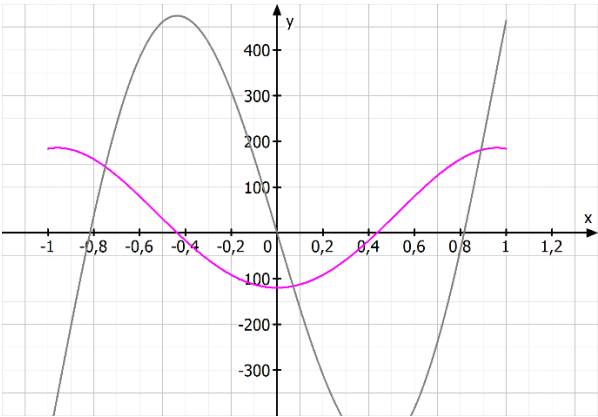
$$\xi \approx 2.6 \cdot 10^{12} \cdot x = 7.299 \cdot 10^{-3}$$

We plot ξ within the range $-1 \leq \xi \leq 1$



$H_0(\xi) H_1(\xi) H_2(\xi) H_3(\xi)$

$H_4(\xi) H_5(\xi)$



$H_6(\xi) H_7(\xi)$

Orthogonality revisited

We got that the Hermite polynomials are orthogonal by a proof working with partial integration.

We proofed for $m \neq n, m < n$:

$$\int_{-\infty}^{\infty} \psi_n(\xi) \cdot \psi_m(\xi) d\xi = 0$$

We did this by deriving $H_m(\xi)$ n -times with the effect:

$$\left(\frac{d}{d\xi}\right)^n H_m(\xi) = 0$$

What we got was:

We examine: Due to $e^{-\xi^2}$ we have:

$$\left[H_m(\xi) \cdot (e^{-\xi^2}) \right]_{-\infty}^{\infty} = 0$$

The integral:

$$\int_{-\infty}^{\infty} \left(\left(\frac{d}{d\xi}\right)^n (H_m(\xi)) \right) \cdot (e^{-\xi^2})$$

But what would happen if we derived m -times instead of n -times?

We would get:

$$\left[H_m(\xi) \cdot \left(\left(\frac{d}{d\xi}\right)^{(n-m)} e^{-\xi^2} \right) \right]_{-\infty}^{\infty} + \int_{-\infty}^{\infty} \left(\left(\frac{d}{d\xi}\right)^m (H_m(\xi)) \right) \cdot \left(\left(\frac{d}{d\xi}\right)^{(n-m)} e^{-\xi^2} \right) d\xi$$

The first part goes to zero due to $e^{-\xi^2}$:

$$\left[H_m(\xi) \cdot \left(\left(\frac{d}{d\xi}\right)^{(n-m)} e^{-\xi^2} \right) \right]_{-\infty}^{\infty} = 0$$

The second part:

$$\int_{-\infty}^{\infty} \left(\left(\frac{d}{d\xi}\right)^m (H_m(\xi)) \right) \cdot \left(\left(\frac{d}{d\xi}\right)^{(n-m)} e^{-\xi^2} \right) d\xi$$

First multiplicand according to (3) gives:

$$\left(\frac{d}{d\xi}\right)^m (H_m(\xi)) = m! \cdot 2^m$$

With $d = n - m$ and (1a) the second multiplicand gives:

$$\left(\frac{d}{d\xi}\right)^d e^{-\xi^2} = H_d(\xi) \cdot (-1)^d \cdot e^{-\xi^2}$$

We combine:

$$\int_{-\infty}^{\infty} \left(\left(\frac{d}{d\xi} \right)^m (H_m(\xi)) \right) \cdot \left(\left(\frac{d}{d\xi} \right)^{(n-m)} e^{-\xi^2} \right) d\xi \rightarrow$$

$$(-1)^d \cdot m! \cdot 2^m \cdot \int_{-\infty}^{\infty} H_d(\xi) \cdot e^{-\xi^2} d\xi$$

We need to calculate $\int_{-\infty}^{\infty} H_d(\xi) \cdot e^{-\xi^2} d\xi$.

For odd Hermite polynomials the integral will give zero because:

$$\int_{-\infty}^{\infty} \xi \cdot e^{-\xi^2} d\xi = 0, \int_{-\infty}^{\infty} \xi^3 \cdot e^{-\xi^2} d\xi = 0, \int_{-\infty}^{\infty} \xi^{2n+1} \cdot e^{-\xi^2} d\xi = 0$$

We take even Hermite polynomials:

$$H_0(\xi) = 1$$

$$H_2(\xi) = 4 \cdot \xi^2 - 2$$

$$H_4(\xi) = 16 \cdot \xi^4 - 48 \cdot \xi^2 + 12$$

$$H_6(\xi) = 64 \cdot \xi^6 - 480 \cdot \xi^4 + 720 \cdot \xi^2 - 120$$

$$H_8(\xi) = 256 \cdot \xi^8 - 3584 \cdot \xi^6 + 13440 \cdot \xi^4 - 13440 \cdot \xi^2 + 1680$$

$$H_{10}(\xi) = 1024 \cdot \xi^{10} - 23040 \cdot \xi^8 + 161280 \cdot \xi^6 - 403200 \cdot \xi^4 + 302400 \cdot \xi^2 - 30240$$

We use:

$$\int_{-\infty}^{\infty} \xi^{2k} \cdot e^{-\xi^2} d\xi = \frac{(2k-1)!!}{2^k} \cdot \sqrt{\pi}$$

We get for $H_2(\xi)$:

$$\int_{-\infty}^{\infty} (4 \cdot \xi^2 - 2) \cdot e^{-\xi^2} d\xi = 4 \cdot \int_{-\infty}^{\infty} \xi^2 d\xi - 2 \cdot \int_{-\infty}^{\infty} e^{-\xi^2} d\xi =$$

$$2 \cdot \sqrt{\pi} - 2 \cdot \sqrt{\pi} = 0$$

We get for $H_4(\xi)$:

$$\int_{-\infty}^{\infty} (16 \cdot \xi^4 - 48 \cdot \xi^2 + 12) \cdot e^{-\xi^2} d\xi =$$

$$16 \cdot \int_{-\infty}^{\infty} \xi^4 \cdot e^{-\xi^2} d\xi - 48 \cdot \int_{-\infty}^{\infty} \xi^2 \cdot e^{-\xi^2} d\xi + 12 \cdot \int_{-\infty}^{\infty} e^{-\xi^2} d\xi =$$

$$16 \cdot \frac{3}{4} \sqrt{\pi} - 48 \cdot \frac{1}{2} \cdot \sqrt{\pi} + 12 \cdot \sqrt{\pi} = 0$$

This cannot be by chance; we check this by help of excel up to $H_{12}(\xi)$:

Hermite polynomials

k	n	$((2k-1)!)/2^k$	H2	H4	H6	H8	H10	H12	H14	H16	H18	H20		
0	0	1	-2	12	-120	1680	-30240	665280	-17297280	518918400	-17643225600	6,70443E+11	x^0	
1	2	0,5	4	-48	720	-13440	302400	-7983360	242161920	-8302694400	3,17578E+11	-1,34089E+13	x^2	
2	4	0,75		16	-480	13440	-403200	13305600	-484323840	19372953600	-8,46875E+11	4,02266E+13	x^4	
3	6	1,875			64	-3584	161280	-7096320	322882560	-15498362880	7,90417E+11	-4,29083E+13	x^6	
4	8	6,5625				256	-23040	1520640	-92252160	5535129600	-3,3875E+11	2,14542E+13	x^8	
5	10	29,53125					1024	-135168	12300288	-984023040	75277762560	-5,72111E+12	x^{10}	
6	12	162,421875						4096	-745472	89456640	-9124577280	8,66835E+11	x^{12}	
7	14	1055,742188							16384	-3932160	601620480	-76205260800	x^{14}	
8	16	7918,066406								65536	-20054016	3810263040	x^{16}	
9	18	67303,56445									262144	-99614720	x^{18}	
10	20	639383,8623	We build the sum "-2*1+4*0.5" etc.										1048576	x^{20}
11	22	6713530,554	0	0	0	0	0	0	0	0	0	0	0	

What we can assume (and should be proofed) is that $\int_{-\infty}^{\infty} H_n(\xi) \cdot e^{-\xi^2} d\xi$ is zero for every Hermite polynomial with exception of $H_0(\xi)$.