

We analyze the model of a general two-level system in a magnetic field using the density matrix formalism instead of pure state vectors.

You may find more information at:

https://ocw.mit.edu/courses/8-05-quantum-physics-ii-fall-2013/7c376a3b5b27e5ea586a7e37b85a05d3_MIT8_05F13_Chap_07.pdf

[https://phys.libretexts.org/Bookshelves/Quantum_Mechanics/Essential_Graduate_Physics_-_Quantum_Mechanics_\(Likharev\)/05%3A_Some_Exactly_Solvable_Problems/5.01%3A_Two-level_Systems](https://phys.libretexts.org/Bookshelves/Quantum_Mechanics/Essential_Graduate_Physics_-_Quantum_Mechanics_(Likharev)/05%3A_Some_Exactly_Solvable_Problems/5.01%3A_Two-level_Systems)

Hope I can help you with learning quantum mechanics.

Recall the Hamiltonian:

$$\hat{H} = -\vec{\mu} \cdot \vec{B} = \frac{1}{2} \hbar \begin{pmatrix} \omega_0 & \omega_R e^{-i\omega t} \\ \omega_R e^{i\omega t} & -\omega_0 \end{pmatrix}$$

We can express by help of the Pauli matrix basis:

$$\hat{H} = \frac{1}{2} \hbar \begin{pmatrix} \omega_0 & V_1 - iV_2 \\ V_1 + iV_2 & -\omega_0 \end{pmatrix} = \frac{1}{2} \hbar (V_1 \hat{\sigma}_x + V_2 \hat{\sigma}_y + \omega_0 \hat{\sigma}_z)$$

Note: $\hat{\sigma}_x, \hat{\sigma}_y, \hat{\sigma}_z$ are the Pauli matrices:

$$\hat{\sigma}_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \hat{\sigma}_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \hat{\sigma}_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

We calculate the coefficient matrices.

$$\begin{pmatrix} 0 & V_1 \\ V_1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & -iV_2 \\ iV_2 & 0 \end{pmatrix} = \begin{pmatrix} 0 & \omega_R e^{-i\omega t} \\ \omega_R e^{i\omega t} & 0 \end{pmatrix}$$

We have two equations:

$V_1 - iV_2 = \omega_R e^{-i\omega t}$	$V_1 + iV_2 = \omega_R e^{i\omega t}$
$V_1 = iV_2 + \omega_R e^{-i\omega t}$	
	$iV_2 + \omega_R e^{-i\omega t} + iV_2 = \omega_R e^{i\omega t}$ $2iV_2 = \omega_R (e^{i\omega t} - e^{-i\omega t})$ $V_2 = \frac{1}{2i} \omega_R (e^{i\omega t} - e^{-i\omega t})$ <p>Can be noted as:</p> $V_2 = \omega_R \sin(\omega t)$
$V_1 = i \frac{1}{2i} \omega_R (e^{i\omega t} - e^{-i\omega t}) + \omega_R e^{-i\omega t}$ $V_1 = \frac{1}{2} \omega_R (e^{i\omega t} - e^{-i\omega t} + 2e^{-i\omega t})$ $V_1 = \frac{1}{2} \omega_R (e^{i\omega t} + e^{-i\omega t})$ <p>Can be noted as:</p> $V_1 = \omega_R \cos(\omega t)$	

Equation of motion for Bloch vector

We take the Pauli matrix representation of the general two-level system Hamiltonian:

$$\hat{H} = \frac{1}{2} \hbar (V_1 \hat{\sigma}_x + V_2 \hat{\sigma}_y + \omega_0 \hat{\sigma}_z)$$

We get an equation of motion for the Bloch vector, evolving under this Hamiltonian, by computing the evolution of an arbitrary general two-level system density matrix state $\vec{\rho}$, represented using the Pauli matrices:

$$\vec{\rho} = \frac{Id + \vec{r} \cdot \vec{\sigma}}{2} = \frac{1}{2} (Id + r_1 \hat{\sigma}_x + r_2 \hat{\sigma}_y + r_3 \hat{\sigma}_z) =$$

$$\frac{1}{2} \begin{pmatrix} 1 + r_3 & r_1 - ir_2 \\ r_1 + ir_2 & 1 - r_3 \end{pmatrix}$$

We use:

$$i\hbar \frac{d\vec{\rho}}{dt} = [\hat{H}, \vec{\rho}] \rightarrow \frac{d\vec{\rho}}{dt} = -\frac{i}{\hbar} [\hat{H}, \vec{\rho}]$$

We calculate the equations of motion for the coordinates r_1, r_2, r_3 .

We can do this in two ways.

First by the cross product.

We use:

$$\frac{d\vec{\rho}}{dt} = \vec{\omega} \times \vec{r}, \quad \vec{\omega} = \begin{pmatrix} V_1 \\ V_2 \\ \omega_0 \end{pmatrix}, \quad \vec{r} = \begin{pmatrix} r_1 \\ r_2 \\ r_3 \end{pmatrix}$$

We calculate:

$$\vec{\omega} \times \vec{r} = \begin{pmatrix} V_2 r_3 - \omega_0 r_2 \\ \omega_0 r_1 - V_1 r_3 \\ V_1 r_2 - V_2 r_1 \end{pmatrix}$$

We get the equations of motion for r_1, r_2, r_3 :

$$\dot{r}_1 = V_2 r_3 - \omega_0 r_2$$

$$\dot{r}_2 = \omega_0 r_1 - V_1 r_3$$

$$\dot{r}_3 = V_1 r_2 - V_2 r_1$$

Second by calculating the commutator:

$$\frac{d\vec{\rho}}{dt} = -\frac{i}{\hbar} [\hat{H}, \vec{\rho}] = -\frac{i}{\hbar} (\hat{H}\vec{\rho} - \vec{\rho}\hat{H})$$

We use:

$$\hat{H} = \frac{1}{2} \hbar \begin{pmatrix} \omega_0 & \omega_R e^{-i\omega t} \\ \omega_R e^{i\omega t} & -\omega_0 \end{pmatrix} = \frac{1}{2} \hbar \begin{pmatrix} \omega_0 & V_1 - iV_2 \\ V_1 + iV_2 & -\omega_0 \end{pmatrix}$$

$$\rho = \begin{pmatrix} \rho_{11} & \rho_{12} \\ \rho_{21} & \rho_{22} \end{pmatrix} = \frac{1}{2} \begin{pmatrix} r_0 + r_3 & r_1 - ir_2 \\ r_1 + ir_2 & r_0 - r_3 \end{pmatrix}$$

We get:

$$\begin{aligned} \hat{H}\rho &= \frac{1}{4} \hbar \begin{pmatrix} \omega_0 & V_1 - iV_2 \\ V_1 + iV_2 & -\omega_0 \end{pmatrix} \begin{pmatrix} r_0 + r_3 & r_1 - ir_2 \\ r_1 + ir_2 & r_0 - r_3 \end{pmatrix} = \\ & \frac{1}{4} \hbar \begin{pmatrix} \omega_0(r_0 + r_3) + (V_1 - iV_2)(r_1 + ir_2) & \omega_0(r_1 - ir_2) + (V_1 - iV_2)(r_0 - r_3) \\ (V_1 + iV_2)(r_0 + r_3) - \omega_0(r_1 + ir_2) & (V_1 + iV_2)(r_1 - ir_2) - \omega_0(r_0 - r_3) \end{pmatrix} \end{aligned}$$

$$\begin{aligned} \rho\hat{H} &= \frac{1}{4} \hbar \begin{pmatrix} r_0 + r_3 & r_1 - ir_2 \\ r_1 + ir_2 & r_0 - r_3 \end{pmatrix} \begin{pmatrix} \omega_0 & V_1 - iV_2 \\ V_1 + iV_2 & -\omega_0 \end{pmatrix} = \\ & \frac{1}{4} \hbar \begin{pmatrix} (r_0 + r_3)\omega_0 + (r_1 - ir_2)(V_1 + iV_2) & (r_0 + r_3)(V_1 - iV_2) - (r_1 - ir_2)\omega_0 \\ (r_1 + ir_2)\omega_0 + (r_0 - r_3)(V_1 + iV_2) & (r_1 + ir_2)(V_1 - iV_2) - (r_0 - r_3)\omega_0 \end{pmatrix} \end{aligned}$$

We build the differences $\hat{H}\rho - \rho\hat{H}$:

$\hat{H}\rho_{11} - \rho\hat{H}_{11}$:

$$\begin{aligned} & \omega_0(r_0 + r_3) + (V_1 - iV_2)(r_1 + ir_2) - (r_0 + r_3)\omega_0 - (r_1 - ir_2)(V_1 + iV_2) = \\ & (V_1 - iV_2)(r_1 + ir_2) - (r_1 - ir_2)(V_1 + iV_2) = \\ & V_1r_1 - iV_2r_1 + iV_1r_2 + V_2r_2 - (V_1r_1 + iV_2r_1 - iV_1r_2 + V_2r_2) = \\ & \quad V_1r_1 - iV_2r_1 + iV_1r_2 + V_2r_2 - V_1r_1 - iV_2r_1 + iV_1r_2 - V_2r_2 = \\ & \quad -2iV_2r_1 + 2iV_1r_2 \end{aligned}$$

$\hat{H}\rho_{12} - \rho\hat{H}_{12}$:

$$\begin{aligned} & \omega_0(r_1 - ir_2) + (V_1 - iV_2)(r_0 - r_3) - ((r_0 + r_3)(V_1 - iV_2) - (r_1 - ir_2)\omega_0) = \\ & \omega_0r_1 - \omega_0ir_2 + V_1r_0 - V_1r_3 - iV_2r_0 + iV_2r_3 - (r_0V_1 - r_0iV_2 + r_3V_1 - r_3iV_2 - r_1\omega_0 + ir_2\omega_0) = \\ & \omega_0r_1 - \omega_0ir_2 + V_1r_0 - V_1r_3 - iV_2r_0 + iV_2r_3 - r_0V_1 + r_0iV_2 - r_3V_1 + r_3iV_2 + r_1\omega_0 - ir_2\omega_0 = \\ & \omega_0r_1 - \omega_0ir_2 + V_1r_0 - V_1r_3 - iV_2r_0 + iV_2r_3 - r_0V_1 + r_0iV_2 - r_3V_1 + r_3iV_2 + r_1\omega_0 - ir_2\omega_0 = \\ & \quad 2\omega_0r_1 - 2\omega_0ir_2 - 2V_1r_3 + 2iV_2r_3 \end{aligned}$$

$\hat{H}\rho_{21} - \rho\hat{H}_{21}$:

$$\begin{aligned} & (V_1 + iV_2)(r_0 + r_3) - \omega_0(r_1 + ir_2) - ((r_1 + ir_2)\omega_0 + (r_0 - r_3)(V_1 + iV_2)) = \\ & V_1r_0 + V_1r_3 + iV_2r_0 + iV_2r_3 - \omega_0r_1 - \omega_0ir_2 - (\omega_0r_1 + \omega_0ir_2 + V_1r_0 + iV_2r_0 - V_1r_3 - iV_2r_3) = \\ & \quad V_1r_0 + V_1r_3 + iV_2r_0 + iV_2r_3 - \omega_0r_1 - \omega_0ir_2 - \omega_0r_1 - \omega_0ir_2 - V_1r_0 - iV_2r_0 + V_1r_3 + iV_2r_3 = \\ & \quad 2V_1r_3 + 2iV_2r_3 - 2\omega_0r_1 - 2\omega_0ir_2 \end{aligned}$$

$\hat{H}\rho_{22} - \rho\hat{H}_{22}$:

$$\begin{aligned} & (V_1 + iV_2)(r_1 - ir_2) - \omega_0(r_0 - r_3) - ((r_1 + ir_2)(V_1 - iV_2) - (r_0 - r_3)\omega_0) = \\ & V_1r_1 - V_1ir_2 + iV_2r_1 + V_2r_2 - \omega_0r_0 + \omega_0r_3 - (r_1V_1 - r_1iV_2 + ir_2V_1 + r_2V_2 - \omega_0r_0 + \omega_0r_3) = \\ & \quad V_1r_1 - V_1ir_2 + iV_2r_1 + V_2r_2 - \omega_0r_0 + \omega_0r_3 - r_1V_1 + r_1iV_2 - ir_2V_1 - r_2V_2 + \omega_0r_0 - \omega_0r_3 = \\ & \quad V_1r_1 - V_1ir_2 + iV_2r_1 + V_2r_2 - \omega_0r_0 + \omega_0r_3 - r_1V_1 + r_1iV_2 - ir_2V_1 - r_2V_2 + \omega_0r_0 - \omega_0r_3 = \\ & \quad -2V_1ir_2 + 2iV_2r_1 \end{aligned}$$

We get the complete commutator:

$$\begin{aligned} & \frac{1}{4}\hbar \left(\cdot -\frac{i}{\hbar} \begin{pmatrix} -2iV_2r_1 + 2iV_1r_2 & 2\omega_0r_1 - 2\omega_0ir_2 - 2V_1r_3 + 2iV_2r_3 \\ 2V_1r_3 + 2iV_2r_3 - 2\omega_0r_1 - 2\omega_0ir_2 & -2V_1ir_2 + 2iV_2r_1 \end{pmatrix} \right) \\ & \quad -\frac{i}{2} \begin{pmatrix} -iV_2r_1 + iV_1r_2 & \omega_0r_1 - \omega_0ir_2 - V_1r_3 + iV_2r_3 \\ V_1r_3 + iV_2r_3 - \omega_0r_1 - \omega_0ir_2 & -V_1ir_2 + iV_2r_1 \end{pmatrix} \\ & \quad \frac{1}{2} \begin{pmatrix} -V_2r_1 + V_1r_2 & -i\omega_0r_1 - \omega_0r_2 + iV_1r_3 + V_2r_3 \\ -iV_1r_3 + V_2r_3 + i\omega_0r_1 - \omega_0r_2 & -V_1r_2 + V_2r_1 \end{pmatrix} \\ & \quad \frac{1}{2} \begin{pmatrix} V_1r_2 - V_2r_1 & i(V_1r_3 - \omega_0r_1) + V_2r_3 - \omega_0r_2 \\ i(\omega_0r_1 - V_1r_3) + V_2r_3 - \omega_0r_2 & V_2r_1 - V_1r_2 \end{pmatrix} \end{aligned}$$

We use the equation of motion:

$$\vec{\rho} = \frac{Id + \vec{r} \cdot \vec{\sigma}}{2} = \frac{1}{2}(Id + r_1 \hat{\sigma}_x + r_2 \hat{\sigma}_y + r_3 \hat{\sigma}_z) \rightarrow$$

$$\frac{d\vec{\rho}}{dt} = \frac{1}{2}(\dot{r}_1 \hat{\sigma}_x + \dot{r}_2 \hat{\sigma}_y + \dot{r}_3 \hat{\sigma}_z)$$

We use the Pauli matrices:

$$\hat{\sigma}_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad \hat{\sigma}_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \quad \hat{\sigma}_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

We expand:

$$\dot{r}_1 \hat{\sigma}_x + \dot{r}_2 \hat{\sigma}_y + \dot{r}_3 \hat{\sigma}_z = \begin{pmatrix} 0 & \dot{r}_1 \\ \dot{r}_1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & -i\dot{r}_2 \\ i\dot{r}_2 & 0 \end{pmatrix} + \begin{pmatrix} \dot{r}_3 & 0 \\ 0 & -\dot{r}_3 \end{pmatrix} =$$

$$\begin{pmatrix} \dot{r}_3 & \dot{r}_1 - i\dot{r}_2 \\ \dot{r}_1 + i\dot{r}_2 & -\dot{r}_3 \end{pmatrix}$$

We compare:

$$\frac{1}{2} \begin{pmatrix} V_1 r_2 - V_2 r_1 & i(V_1 r_3 - \omega_0 r_1) + V_2 r_3 - \omega_0 r_2 \\ i(\omega_0 r_1 - V_1 r_3) + V_2 r_3 - \omega_0 r_2 & V_2 r_1 - V_1 r_2 \end{pmatrix}$$

We get:

$\dot{r}_1 = V_2 r_3 - \omega_0 r_2$	$\dot{r}_2 = \omega_0 r_1 - V_1 r_3$	$\dot{r}_3 = V_1 r_2 - V_2 r_1$
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This fits with the cross product:

$$\vec{\omega} \times \vec{r} = \begin{pmatrix} V_2 r_3 - \omega_0 r_2 \\ \omega_0 r_1 - V_1 r_3 \\ V_1 r_2 - V_2 r_1 \end{pmatrix}$$