

One of the problems representing spin is that in three-dimensional space we have three directions: x, y, and z. However, the spin state space works with complex-valued vector spaces. A two-dimensional complex vector space corresponds to a four-dimensional real vector space, which we can no longer easily visualize.

We express the possible spin states in the  $|0\rangle/|1\rangle$  basis. This is the orientation of the spin along the z-axis.

Spin states along the z-axis:	Spin states along the x-axis:	Spin along the y-axis:
$ 0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$  $ 1\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$	$ x+\rangle = \frac{1}{\sqrt{2}}( 0\rangle +  1\rangle)$ $ x+\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$ $ x-\rangle = \frac{1}{\sqrt{2}}( 0\rangle -  1\rangle)$ $ x-\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$	$ y+\rangle = \frac{1}{\sqrt{2}}( 0\rangle + i 1\rangle)$ $ y+\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ i \end{pmatrix}$ $ y-\rangle = \frac{1}{\sqrt{2}}( 0\rangle - i 1\rangle)$ $ y-\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -i \end{pmatrix}$

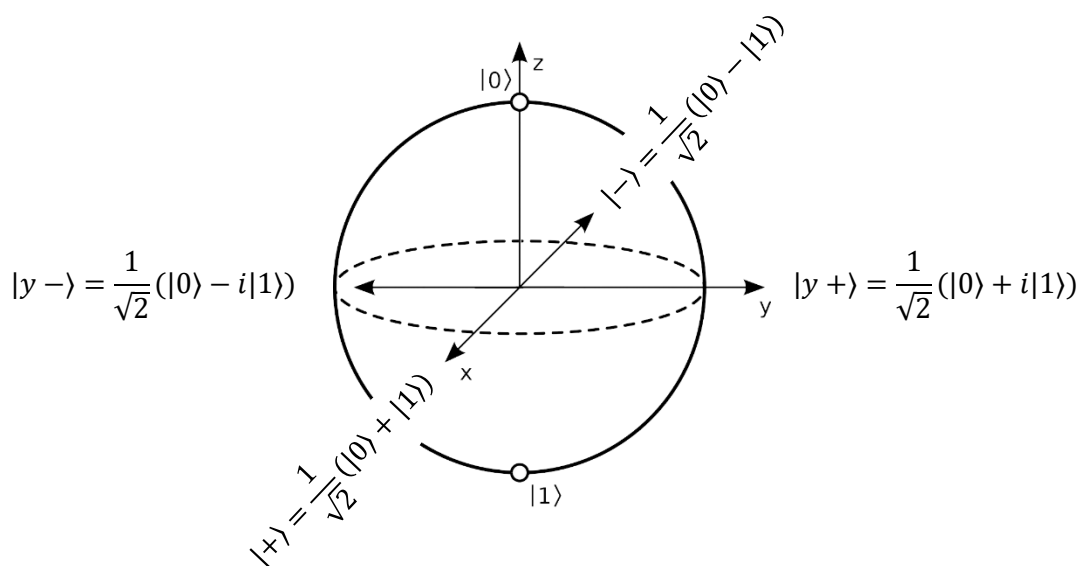
Note:  $|x+\rangle$  usually written as  $|+\rangle$ ,  $|x-\rangle$  usually written as  $|-\rangle$

Note:  $|y+\rangle$  often written as  $|i+\rangle$ ,  $|y-\rangle$  often written as  $|i-\rangle$

The orientation of spin is visualized by using the Bloch-Sphere.

A spin orientated spherical in the z-direction is written as  $|0\rangle$  and corresponds to the complex-valued vector  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$ . A spin orientated spherical in the  $-z$ -direction is written as  $|1\rangle$  and corresponds to the complex-valued vector  $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$ . Please note that  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$  and  $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$  are orthogonal vectors, whereas the spherical direction z and  $-z$  are not.

Using the  $|0\rangle/|1\rangle$ -basis we assign:



Please note that the spherical directions  $x$  and  $-x$  as well as  $y$  and  $-y$  are not orthogonal, the states  $|+\rangle$ ,  $|-\rangle$  and  $|y+\rangle$ ,  $|y-\rangle$  are orthogonal.

An arbitrary state  $|\psi\rangle$  can be expressed in the  $|0\rangle/|1\rangle$ -basis by first rotating it by the angle  $\theta$  with respect to the  $z$ -axis and then by the angle  $\varphi$  with respect to the  $x$ -axis.

We have to deal with the spherical/complex problematic again.

Rotating a spherical vector by  $90^\circ$  resp. by  $\frac{\pi}{2}$  produces an orthogonal spherical vector.

Rotating a complex state by  $180^\circ$  resp. by  $\pi$  produces an orthogonal complex state.

We show this by an example, setting  $\theta = \frac{\pi}{2}$  and  $\varphi = 0$ . Obviously, the state  $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$  will change to the state  $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$

In contrast, if we take the spherical vector  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$  and rotate it by  $\frac{\pi}{2}$  it will change into the vector  $-\begin{pmatrix} 0 \\ 1 \end{pmatrix}$  and this is an orthogonal direction spherical.

Rotating is performed by using unitary matrices. An invertible complex valued square matrix  $U$  is unitary if its conjugate transpose is also its inverse:

$$U^\dagger U = U U^\dagger = I$$

Rotation preserves the length (norm) of a vector. In quantum mechanics the length of a vector corresponds to its probability amplitude, so a rotation doesn't change it.

Note: The rows and columns of a unitary matrix form an orthonormal basis of  $\mathbb{C}^n$ .

Rotation around one of the primary axes is performed by help of the Pauli matrices:

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

The Pauli matrices are unitary:

$$\sigma_x^\dagger = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma_y^\dagger = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma_z^\dagger = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

We check unitarity:

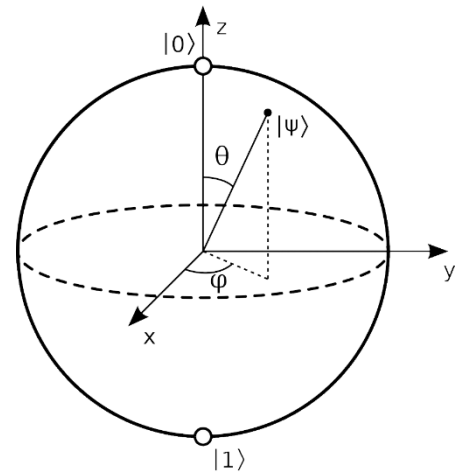
$$\sigma_x^\dagger \sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad \{1\}$$

$$\sigma_y^\dagger \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad \{2\}$$

$$\sigma_z^\dagger \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \quad \{3\}$$

We apply rotation around the  $y$ -axis to the state vector  $|0\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$ :

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ i \end{pmatrix} = -i \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \{4\}$$



We call the factor  $-i$  a global phase. A global phase has no effect to quantum mechanical measurements and can be ignored in most cases.

We rewrite  $-i = e^{i\frac{3}{2}\pi}$  and get:

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = e^{i\frac{3}{2}\pi} \cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

Obviously, the rotation angle is  $\pi$  resp.  $180^\circ$ . Rotating twice by use of the Pauli matrices restores the original state:

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} e^{i\frac{3}{2}\pi} \cdot \begin{pmatrix} 0 \\ 1 \end{pmatrix} = e^{i\frac{3}{2}\pi} \cdot -i \begin{pmatrix} 1 \\ 0 \end{pmatrix} = e^{i\frac{6}{2}\pi} \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \{5\}$$

Note:  $e^{i\frac{6}{2}\pi} = e^{i2\pi} = 1$

The vector  $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$  should rotate to the vector  $|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$ . We check this:

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = -i \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} \quad \{6\}$$

Again, we use that  $-i = e^{i\frac{3}{2}\pi}$  and get:

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = e^{i\frac{3}{2}\pi} \cdot \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

We rotate again:

$$\begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} e^{i\frac{3}{2}\pi} \cdot \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = e^{i\frac{3}{2}\pi} \cdot i \cdot \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \quad \{7\}$$

$$e^{i\frac{3}{2}\pi} \cdot e^{i\frac{1}{2}\pi} \cdot \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix}$$

Besides working with explicit representation of vectors we can work on the conceptual level.

The calculations above, {1} til {7} give:

$$\{1\} \sigma_x^\dagger \sigma_x = Id, \{2\} \sigma_y^\dagger \sigma_y = Id, \{3\} \sigma_z^\dagger \sigma_z = Id$$

$$\{4\} \sigma_y |0\rangle = |1\rangle$$

$$\{5\} \sigma_y |1\rangle = |0\rangle$$

$$\{6\} \sigma_y |+\rangle = |-\rangle$$

$$\{7\} \sigma_y |-\rangle = |+\rangle$$

So far, we can rotate the Bloch sphere around the axes  $x, y, z$  by  $\pi$  resp.  $180^\circ$ .

If we want to interchange two points on the Bloch sphere, we can build this rotation with a combination of a maximum of three rotations around the axes  $x, y, z$ .

In order to rotate the Bloch sphere round the  $x/y/z$ -axis by an angle of  $\theta$  we apply the unitary matrix:

$$e^{-i\frac{\theta}{2}\sigma_x} = \begin{pmatrix} \cos\left(\frac{\theta}{2}\right) & -i \cdot \sin\left(\frac{\theta}{2}\right) \\ -i \cdot \sin\left(\frac{\theta}{2}\right) & \cos\left(\frac{\theta}{2}\right) \end{pmatrix}$$

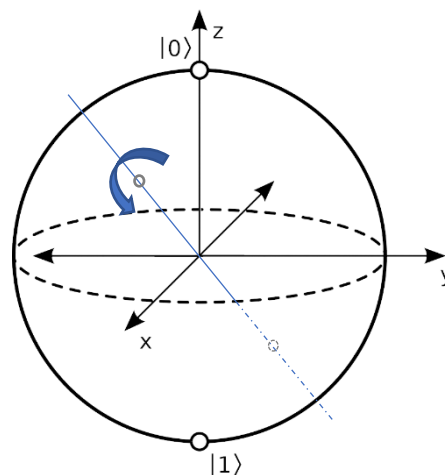
$$e^{-i\frac{\theta}{2}\sigma_y} = \begin{pmatrix} \cos\left(\frac{\theta}{2}\right) & -\sin\left(\frac{\theta}{2}\right) \\ \sin\left(\frac{\theta}{2}\right) & \cos\left(\frac{\theta}{2}\right) \end{pmatrix}$$

$$e^{-i\frac{\theta}{2}\sigma_z} = \begin{pmatrix} e^{-i\frac{\theta}{2}} & 0 \\ 0 & e^{i\frac{\theta}{2}} \end{pmatrix} = e^{-i\frac{\theta}{2}} \begin{pmatrix} 1 & 0 \\ 0 & e^{i\theta} \end{pmatrix} \approx \begin{pmatrix} 1 & 0 \\ 0 & e^{i\theta} \end{pmatrix}$$

Note: for the proof you may refer to

“exponentiation\_of\_pauli\_matrices” on this website within the part of quantum mechanics or [https://www.researchgate.net/publication/335654130 Rotations on the Bloch Sphere](https://www.researchgate.net/publication/335654130_Rotations_on_the_Bloch_Sphere).

The Hadamard matrix  $H$  rotates around the  $(x + z)$ -axis by  $\pi$  resp.  $180^\circ$ :



It flips  $x$ - and  $z$ -axis:

$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \left( \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \begin{pmatrix} 0 \\ 1 \end{pmatrix} \right)$	$H 0\rangle =  +\rangle$
$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 2 \\ 0 \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$	$H +\rangle =  0\rangle$
$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \frac{1}{\sqrt{2}} \left( \begin{pmatrix} 1 \\ 0 \end{pmatrix} - \begin{pmatrix} 0 \\ 1 \end{pmatrix} \right)$	$H 1\rangle =  -\rangle$
$\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 0 \\ 2 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$	$H -\rangle =  1\rangle$